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Correlations between Global Positioning System and U.S. Naval Observatory Master Clock Time

by Thomas B. Bahder

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Abstract

The U.S. Naval Observatory Master Clock is used to steer Global Positioning System (GPS) time. Time-transfer data consisting of the difference between the Master Clock time and the GPS time were recorded from all satellites in the GPS constellation over a time period covering 10 October to 12 December 1995. A Fourier analysis of these data shows a distinct peak in the Fourier spectrum, corresponding approximately to a one-day period. For a more accurate determination of this period, correlations are computed between successive days of data. An average of 25 correlation functions shows a correlation equal to 0.52 at a delay time of 23 hr 56 min (which corresponds to twice the average GPS satellite period). This correlation indicates that GPS time, as measured by the U.S. Naval Observatory, is periodic with respect to the Master Clock, with a period of 23 hr and 56 min. An autocorrelation of a five-day segment of data indicates that these correlations persist for four successive days.

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1. Introduction

The Global Positioning System (GPS) is made up of a constellation of 24 satellites in four orbital planes about the earth [1–4]. Since full operational capability of the GPS was announced, applications of the GPS have grown exponentially as the system’s accuracy, global coverage, and reliability were recognized. However, military [5] and civilian [6] GPS users are already requesting position and time-transfer accuracy beyond that of the system’s original design requirements [2]. To improve GPS accuracy, we must revisit the basic underpinnings of the system, including basic physics, time transfer/steering techniques, and current algorithms used to operate the GPS. In this work, I investigate correlations in time difference between the official Department of Defense (DoD) reference standard (the U.S. Naval Observatory (USNO) Master Clock) and GPS time, as determined from single satellite time transfer [7–11].

The GPS is based on measurement of the pseudorange, which is the phase of the pseudorandom noise (PRN) code that is broadcast from a satellite [1,12]. This phase corresponds to the time of flight of the signal from satellite to receiver (modulo the code period). The measured pseudorange from four satellites can be used to determine the user’s time and position, relative to the known positions of the satellites in the GPS constellation [1]. If all satellite clocks were “perfectly synchronized” with each other [7,10,13–15], and the user’s clock were “perfectly synchronized” with the satellite clocks, the light travel time from three satellites would, in principle, provide sufficient information to determine the user’s position in three-dimensional space [3,4,16]. In practice, however, there is some offset between the user’s clock and each satellite clock. This offset is determined during the measurement process, if pseudorange measurements are made to four or more satellites. Alternatively, if we know the user’s position, and we have the correct time (such as given by the USNO Master Clock) then we can obtain from GPS the time difference between our clock and GPS time by acquiring the navigation message (which contains the satellite ephemeris) and by measuring the pseudorange to one satellite [17]. This is the type of measurement that is routinely made by USNO when GPS time is compared to time on the Master Clock [17].

The GPS epoch is 0000 UT (midnight) on January 6, 1980. GPS time is not adjusted by the addition of leap seconds and is therefore, offset from Coordinated Universal Time (UTC) by an integer number of seconds, plus some tens of nanoseconds. As of 1 January 1996, GPS time is ahead of UTC by 11 s. The relation between GPS time t_{GPS} and the best estimate of UTC(USNO) as obtained from GPS, t_{UTC} , is given by [20,21]

$$t_{UTC} = t_{GPS} - \Delta t_{UTC} \quad (1)$$

where

$$\Delta t_{UTC} = \Delta t_{LS} + A_0 + A_1(t_{GPS} - t_{RT}) \quad (2)$$

where Δt_{LS} is the integer leap seconds time offset between GPS and UTC time, t_{RT} is the reference time for the UTC data, and A_0 and A_1 are constants in the navigation message. The UTC(USNO) time scale is kept within approximately 100 ns of the international time standard, UTC, which is published by the Bureau International des Poids et Mesures (BIPM).

The USNO Master Clock is a hydrogen maser that provides the best real-time estimate of the UTC(USNO) time scale. Hydrogen masers are very stable over short time periods, such as a week. USNO's current Sigma Tau hydrogen masers (STSC Model 2010 (1994)) exhibit a time domain maximum instability of 3.0×10^{-15} for periods of 1000 to 10,000 s. However, the Master Clock is steered to the UTC(USNO) time scale, which itself is based on an ensemble of dozens of cesium clocks and five to ten hydrogen masers. Consequently, the UTC(USNO) time scale is very stable. Its rate does not change by more than about 100 ps per day, from day to day.

GPS time is steered to UTC(USNO). During the last several years, these times have been kept within a few hundred nanoseconds. However, data from single-frequency receivers indicate that there is a diurnal variation in the time difference between GPS and UTC(USNO) times [18,19]. In this work, I investigate the residuals of differences between UTC(USNO) Master Clock time and GPS time, as seen through individual GPS satellites, after correcting each satellite clock by using the broadcast a_{f0} , a_{f1} , and a_{f2} parameters and correcting for ionospheric delay (using an algorithm based on L1 and L2 frequencies) [20].

For DoD purposes, the USNO Master Clock is the best real-time source of time for the time scale USNO(UTC). This time scale is a realization of coordinate time on the geoid, in the rotating Earth-centered, Earth-fixed (ECEF)

frame of reference [10,11]. As the Master Clock moves through inertial space, its (coordinate) time agrees with the coordinate time on hypothetical coordinate clocks with which it instantaneously coincides in the underlying Earth-centered inertial (ECI) frame of reference (i.e., the rate of coordinate clocks in this ECI frame is set by time on the geoid*). The GPS is intended to keep coordinate time in this ECI frame (modulo leap seconds). Specifically, the GPS time that a satellite provides to a user is intended to be the same as that of a (stationary) coordinate clock in the underlying ECI frame, at the position of the user, having the rate of clocks on the geoid. (Ashby [10] gives a detailed discussion of GPS time as coordinate time.) Ideally, if all facets of the GPS are implemented correctly, and time transfer is performed correctly, the observed difference between the two sources of time—the USNO Master Clock and GPS—should contain only random measurement errors, but no systematic differences.

*The ECI frame referred to is not a true inertial frame of reference, as is commonly used in special relativity. This ECI frame is described by a Schwarzschild-like metric with a transformed coordinate time. See Ashby [10].

2. Data Analysis

USNO personnel provided me with data consisting of time differences between UTC(USNO) and GPS times, acquired through the Precise Positioning Service (PPS) via P(Y)-code (encrypted precision code). The data were acquired from all satellites in the GPS constellation [22] during 10 October 1995 to 12 December 1995, which corresponds to modified Julian date (MJD) 50000.00553 to 50063.54009. Pseudorange data were collected by a receiver manufactured by Stanford Telecommunications, Inc. (model 5401C, serial number 021); this dual-frequency (L1 and L2) P(Y)-code receiver is used by USNO to monitor GPS time. USNO made the standard two-frequency correction for ionospheric delay [20]. An algorithm internal to the receiver was used to correct for propagation delay through the troposphere. The broadcast ephemeris was used from each satellite in the receiver's time-transfer computation. Furthermore, two other corrections were made: a Sagnac correction for the Earth rotation, and the relativistic " $e \sin E$ " correction of the signal emission time due to eccentricity of the satellite orbit [20].

Each satellite in turn was tracked, and time-transfer data were collected every 6 s, over a track period of 780 s (13 min) [22]. A least-squares fit to the 6-s data was done (over the 13-min track period) to obtain one data point, which represents a best estimate of the difference between the USNO Master Clock and GPS time, via the individual satellite, at the given MJD (or fraction thereof). The 13-min track period was chosen so that the entire navigation message would be received; the navigation message, transmitted every 12.5 min, includes the latest ionospheric and UTC information. The data consisted of 5623 points total, corresponding to approximately 89 points per day. The root mean square (RMS) deviations for the 13-minute-track data points were on the order of 2 to 10 ns.

Figure 1 is a plot of the original data, consisting of the time difference between the USNO Master Clock and GPS time versus the MJD. Over the 63 days of data, the time difference varied from approximately -40 to $+50$ ns. The long-term variation in the data reflects the complicated response of the GPS to time steering, which includes the Kalman filter process (run at the Master

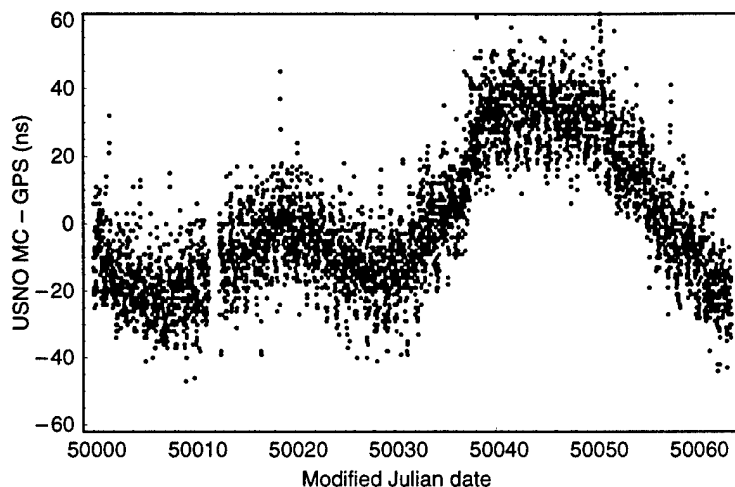


Figure 1: Original data (time difference between USNO Master Clock and GPS time) versus modified Julian date for period MJD 50000.00553 to 50063.54009.

Control Station in Colorado Springs) and the response of other subsystem components. There is a gap in the data during the period MJD 50011.45177 to 50012.46282, due to data acquisition difficulties.

Figure 2 shows a subset of the same data over the time period MJD 50020 to 50028. A salient feature of the data is the scatter of points, on the order of 30 ns peak-to-peak. This scatter is due to a combination of effects [1], including noise in the receiver electronics (~ 0.5 m ≈ 1.5 ns), multipath effects (~ 1.4 m), uncompensated tropospheric delay (~ 0.7 m), uncompensated ionospheric delay (~ 1.2 m), satellite clock errors (~ 2.1 m), and ephemeris errors (~ 2.1 m), where all values are $1-\sigma$ errors in range. The USNO GPS antenna phase center coordinates are believed accurate to approximately 0.5 m. Perhaps the biggest contribution to this scatter in the data arises because the satellites in the constellation do not all have their navigation messages updated at the same time; this variation results in a scatter of values during any given measurement time. Besides this scatter, a diurnal variation is apparent. Figure 3 shows an 11-point average of the data in the time period MJD 50020 to 50028. A diurnal oscillation with a magnitude of 18 to 20 ns peak-to-peak is clearly present. The physical origin of this diurnal variation is not well understood. Workers in the field [18] have proposed a variety of reasons to explain this behavior, including broadcast ephemeris errors, multipath er-

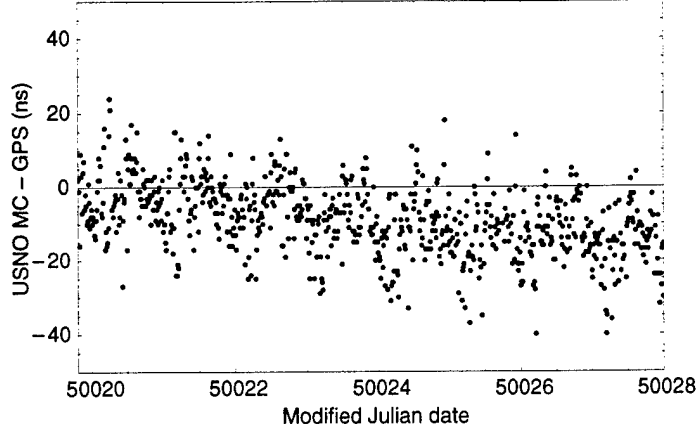


Figure 2: Portion of data from MJD 50020 to 50028, on an expanded scale.

rors, incorrect receiver antenna phase center coordinates, the fundamental accuracy limit of the clocks on the Block II satellite vehicles, poor thermal control of the clock systems (on the ground and in the satellites), inaccuracies in modeling of the ionosphere and troposphere, and the possibility that relativistic effects in the GPS have not been accounted for properly.

2.1 Fourier Transform

I searched for periodicity in the data by performing a Fourier transform, using a fast Fourier transform (FFT) algorithm. The FFT algorithm [23] requires that the data set be uniformly sampled in time; however, the original data set was not uniformly sampled. Therefore, I fit a cubic spline to the original data set and resampled the data at a uniform sampling rate $\Delta = t_{i+1} - t_i = 0.002$ day, where t_i is the time of the i th resampled data point. The sampling rate of the original data varied in the approximate range of 0.009 to 0.013 day, so I have, essentially, lost no information by resampling the data using a smaller sampling interval. For the Fourier transform $H(f)$ of a function $h(t)$, I use the convention

$$H(f_n) = \int_{-\infty}^{+\infty} h(t) e^{2\pi i f_n t} dt \approx \Delta \sum_{k=0}^{N-1} h_k e^{2\pi i f_n t_k} = \Delta \sum_{k=0}^{N-1} h_k e^{2\pi i k n} = \Delta H_n, \quad (3)$$

where $h_k = h(t_k)$. For N data points, there are N Fourier amplitudes at frequencies f_n , where $n = -N/2, \dots, 0, \dots, +N/2$. The finite data set imposes

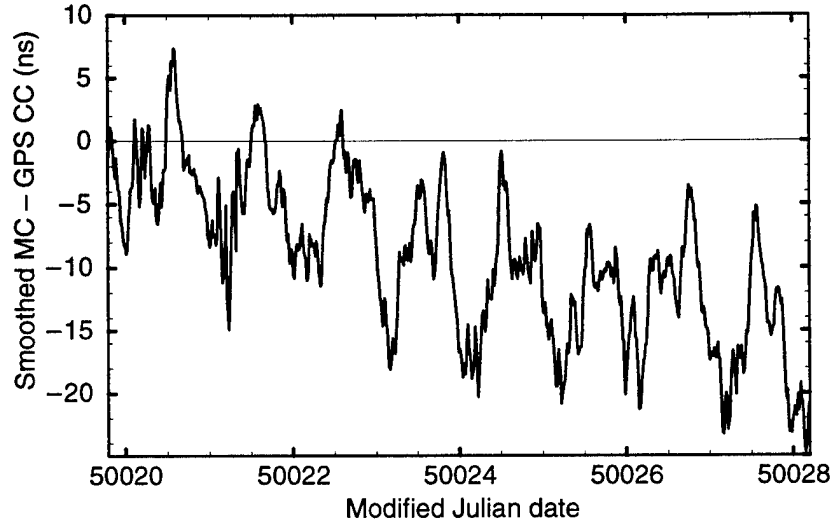


Figure 3: Original data, smoothed by taking average of 11 points, from MJD 50020 to 50028. (A diurnal oscillation of (peak-to-peak) magnitude of 18 to 20 ns is evident.)

the periodicity on the amplitudes $H_{n+N} = H_n$. In particular, the amplitudes $H_{-N/2}$ and $H_{N/2}$ are equal, and not independent. Positive frequencies $0 < f < f_c$, where $f_c = 1/(2\Delta)$ is the Nyquist critical frequency, correspond to discrete f_n for $n = 1, 2, \dots, N/2 - 1$. Furthermore, my data consist of a real function, so the amplitudes for negative frequencies are related to the amplitudes for positive frequencies, $H(-f) = H(f)^*$, or in terms of the discrete amplitudes, $H_{-n} = H_n^*$. Figure 4 shows a plot of the magnitudes of the Fourier transform amplitudes $|H(f_n)|^2$ of the resampled data for the $N/2$ frequencies f_n , where $n = 0, 1, \dots, N/2 - 1$, and $N/2 = 6893$. I cut off the plot at $f = 30 \text{ day}^{-1}$ (a 48-min period), which is just below the approximate Nyquist critical frequency of the original data, above which there is no additional information.*

Figure 4 shows numerous peaks in the Fourier amplitude over a wide range of frequencies. The magnitude of the Fourier amplitude peaks generally decreases with increasing frequency. Most of the periodicity is concentrated at frequencies below $f = 25 \text{ day}^{-1}$, or at periods longer than one hour. At present, I do not understand the large number of strong Fourier amplitudes

*Throughout this work, I use units of day, whose magnitude is given by $1 \text{ day} = 1 \text{ solar day} = 1 \text{ MJD}$.

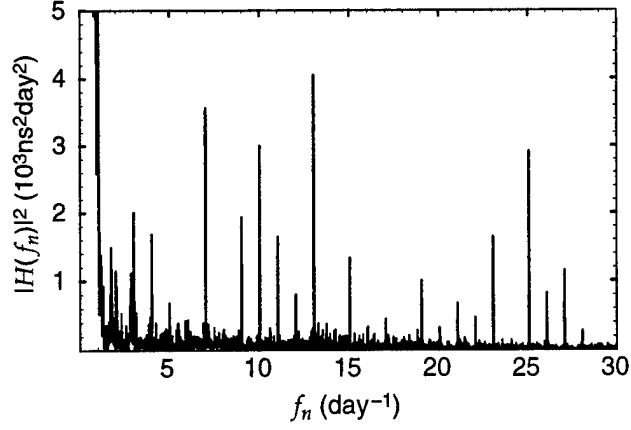


Figure 4: Power spectrum of uniformly resampled data.

in the data. Table 1 lists the prominent peaks in the Fourier amplitudes shown in figure 4. Figure 5 shows the same plot as in figure 4 (on an expanded scale), along with an inset that shows the frequency range $f = 1$ to 8 day^{-1} . A strong peak exists at $f = 0.9915644 \text{ day}^{-1}$, which corresponds to a 1.00851-day period. Two much weaker peaks are present at $f = 1.99887$ and 2.01461 day^{-1} , which correspond to periods of 0.500283 and 0.496374 day, respectively. To the accuracy of the frequency separations in my Fourier analysis, these periodicities correspond to approximate periods of 1.0 and 0.5 day. The 0.5-day period is close to the average period of the GPS satellites, which is $11.9664 \text{ hr} = 0.4986 \text{ day}$.* The discrete frequency spacing prevents determination of a more accurate period than is given by the spacing between discrete frequencies.

Figure 6 shows a further expanded scale near $f = 0$. The strongest peak at $f = 0.0157391 \text{ day}^{-1}$ corresponds to the fundamental time period of the data set, $1/f = 63.536 \text{ day}$. The zero frequency peak represents an overall constant offset of the USNO Master Clock, with respect to GPS time.

*The average GPS satellite period 11.9664 hr is an average of the periods of the 25 satellites; data provided by J. Toth and B. Winn of The Aerospace Corporation.

Table 1: Frequencies, corresponding periods, and absolute squares of Fourier amplitudes shown for prominent peaks in figures 3–5.

f_n (day ⁻¹)	Period (day)	$ H(f_n) ^2$ (10 ³ ns ² day ²)
0	—	14.4223
0.0157391	63.536	808.044
0.125913	7.942	36.1935
0.220348	4.53829	41.3523
0.582347	1.71719	21.8435
0.629564	1.5884	19.4908
0.818434	1.22185	10.2241
0.897129	1.11467	8.31069
0.991564	1.00851	25.1686
1.74704	0.572396	1.48858
1.99887	0.500283	1.14948
2.01461	0.496375	1.03249
2.86452	0.349099	1.11251
3.00617	0.332649	2.00817
4.01347	0.249161	1.68535
5.00504	0.199799	0.682596
7.01964	0.142457	3.56584
9.01851	0.110883	1.93372
10.0258	0.0997425	3.00274
11.0331	0.0906362	1.64622
12.0404	0.0830536	0.805485
13.032	0.0767343	4.05422
15.0466	0.0664603	1.33505
19.0601	0.0524657	1.00523
23.0578	0.0433693	1.64400
25.0724	0.0398845	2.92053
26.0797	0.038344	0.821129
27.0713	0.0369395	1.14859

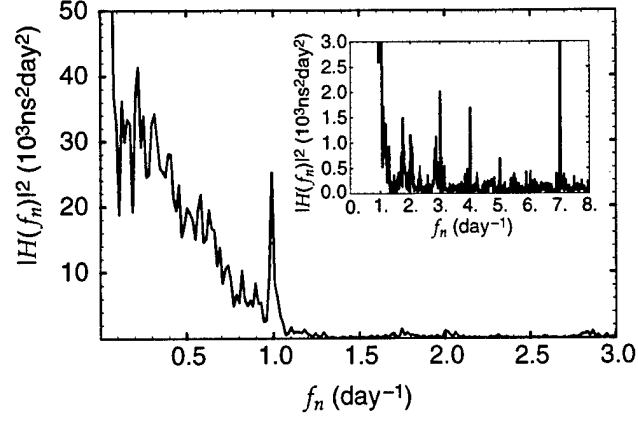


Figure 5: Power spectrum on expanded scale in low-frequency region $f_n = 0$ to 3 day^{-1} . (Note: strong peak is seen at $f_n = 0.991564 \text{ day}^{-1}$. Inset shows same spectrum for intermediate frequency region $f = 1$ to 8 day^{-1} .)

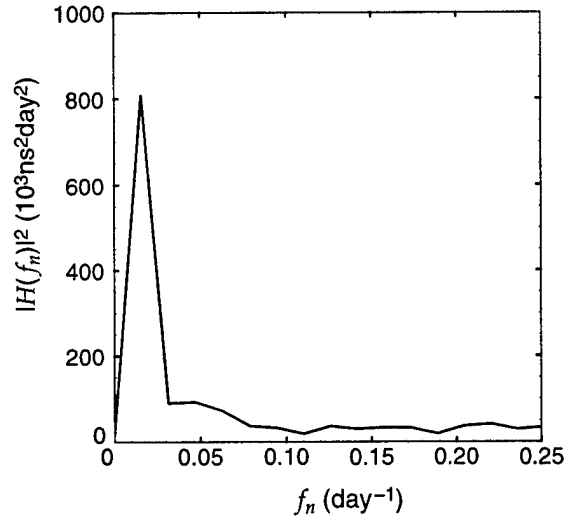


Figure 6: Power spectrum near zero frequency. (Discrete Fourier amplitudes are connected by straight lines.)

2.2 Correlation Function

In order to determine the periodicity of GPS time more accurately, I computed correlations between successive days of the data. If such correlations exist, they may be due to systematic physical effects, which may warrant further exploration. I designate the original data—that is, the USNO Master Clock time minus the GPS time (plotted in fig. 1)—by the function $f(\tilde{m}_i)$, where \tilde{m}_i are the MJD observation times. In terms of a time variable t , these data are given by a function $h(t)$, defined by

$$\begin{aligned} f(\tilde{m}) &= h(t) \\ &= h(t_0 + (\tilde{m} + \kappa_0)\Delta t), \end{aligned} \quad (4)$$

where the time t and MJD \tilde{m} are related by

$$\begin{aligned} t &= t_0 + \tilde{j} \Delta t \\ &= t_0 + (\tilde{m} + \kappa_0) \Delta t, \end{aligned} \quad (5)$$

where \tilde{j} is the Julian date [24], t_0 is the Julian date epoch (Greenwich mean noon on 1 January 4713 BC), $\kappa_0 = 2400000.5$, and $\Delta t = 1.0$ day. To divide the data into segments of length T and correlate the segments with one another, I introduce the function

$$\begin{aligned} h_m(t) &= h(t_0 + (\tilde{m} + \kappa_0) \Delta t) [\theta(\bar{t}) - \theta(\bar{t} - T)] \\ &= h(t_0 + (m + \kappa_0) \Delta t + \bar{t}) [\theta(\bar{t}) - \theta(\bar{t} - T)], \end{aligned} \quad (6)$$

where m is the integer part of the MJD \tilde{m} that labels the starting time of the data segment, \bar{t} is a shifted time variable such that $\bar{t} = 0$ at the start of the data segment, and $\theta(t) = 1$ for $t \geq 0$ and $\theta(t) = 0$ for $t < 0$. (Note that the function $h_m(t)$ is nonzero only over a time interval of length T , which starts at integer MJD m .)

For any given day, the mean of the data $f(\tilde{m}_i)$ is not zero. In order to define a normalized correlation function, I define a new function that has zero mean, $e_n(t)$, by subtracting the mean of $h_m(t)$ over the segment of length T , so that

$$e_n(t) = h_n(t) - \langle h_n(t) \rangle, \quad (7)$$

where the mean for the data starting on integer MJD n is given by

$$\langle h_n(t) \rangle = \frac{1}{T} \int_0^T h_n(t) dt. \quad (8)$$

The functions $e_n(t)$ represent the difference data $f(\tilde{m}_i)$ starting on integer MJD n , with the mean for that data segment subtracted.

I now define the normalized correlation function between two data segments of length T , starting on integer MJD m and MJD n , respectively, by

$$g_{m,n}(\tau) = \frac{1}{(N_m N_n)^{\frac{1}{2}}} \int_0^T e_m(t) e_n(t - \tau) dt, \quad (9)$$

where the normalization constants N_i are given by

$$N_i = \int_0^T e_i^2(t) dt. \quad (10)$$

The correlation function in equation (9) is dimensionless and is defined so that, if the time differences $f(\tilde{m}_i)$ are identical for two data segments, $e_m(t) = e_n(t)$, then $g_{m,n}(0) = 1$. An example of the behavior of $g_{m,n}(\tau)$ for a function that is perfectly correlated from one day to another is shown in figure 7, where I calculated the autocorrelation function, $g_{n,n}(\tau)$, of a 1-day segment of data, starting at MJD $n = 50023$. As $\tau \rightarrow 0$, the autocorrelation function $g_{n,n}(\tau) \rightarrow 1$. For increasing values of τ , $g_{n,n}(\tau)$ decreases for two reasons. First, the data are not correlated at times $\tau > 0$, which leads to a sharp decay of $g_{n,n}(\tau)$. The second reason that $g_{n,n}(\tau)$ decreases for increasing values of τ is that the functions $e_n(t)$ and $e_n(t - \tau)$ have a decreasing overlap for increasing τ . The finite support (finite nonzero domain) of the function $e_n(t)$ leads to less overlap between $e_n(t)$ and $e_n(t - \tau)$ for $\tau > 0$ than for $\tau = 0$, and leads to the decay of the correlation function $g_{n,n}(\tau)$ for increasing τ . This effect in $g_{n,n+1}(\tau)$ is not significant, since (for example) for two constant functions with finite support, $g_{n,n+1}(\tau)$ would decay approximately linearly with τ , whereas I am interested in sharp peaks in the correlation function (such as the peak near $\tau = 0$ in fig. 7).

From the data, I computed the function $g_{n,n+1}(\tau)$, which is a correlation function between two successive days of the data (that is, I used data segments where $T = 1$ day). When the data on two successive days are completely uncorrelated, we would expect $g_{n,n+1}(\tau)$ for $\tau \sim 0$ to be small. Figures 8 to 10 show the computed correlation functions $g_{n,n+1}(\tau)$ for 25 successive pairs of days in the data. Inspection of each correlation function shows that there is a large systematic correlation near $\tau = 0$ in all cases. The time $\tau = 0$ corresponds to correlations with a period of 24 hr. In principle, it is possible

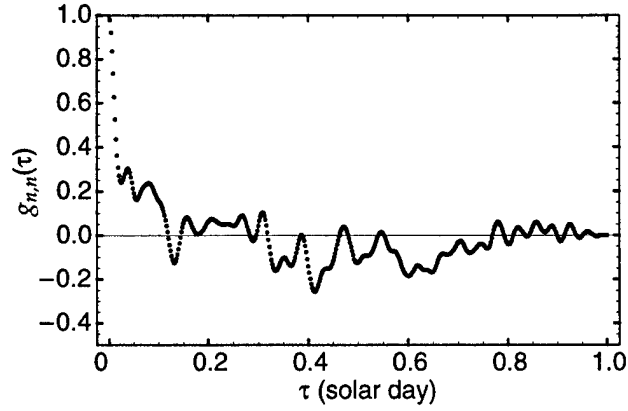


Figure 7: Autocorrelation function, $g_{n,n}(\tau)$, for $T = 1$ day, from MJD $n = 50023$.

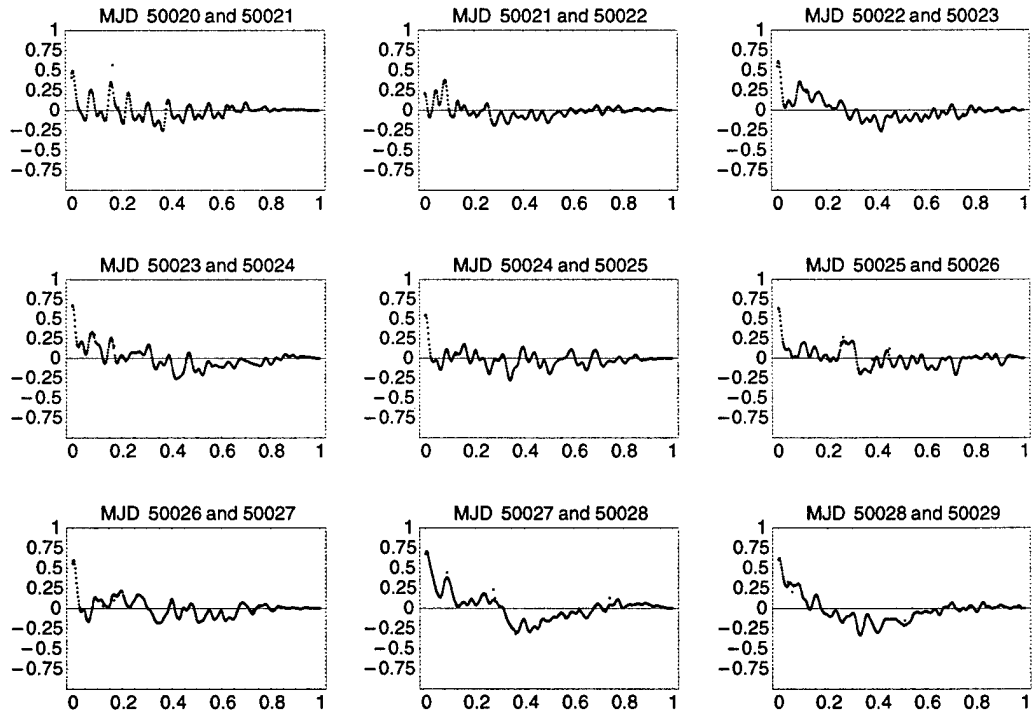


Figure 8: Correlation functions, $g_{n,n+1}(\tau)$, for $T = 1$ day, from MJD $n = 50020$ through 50028.

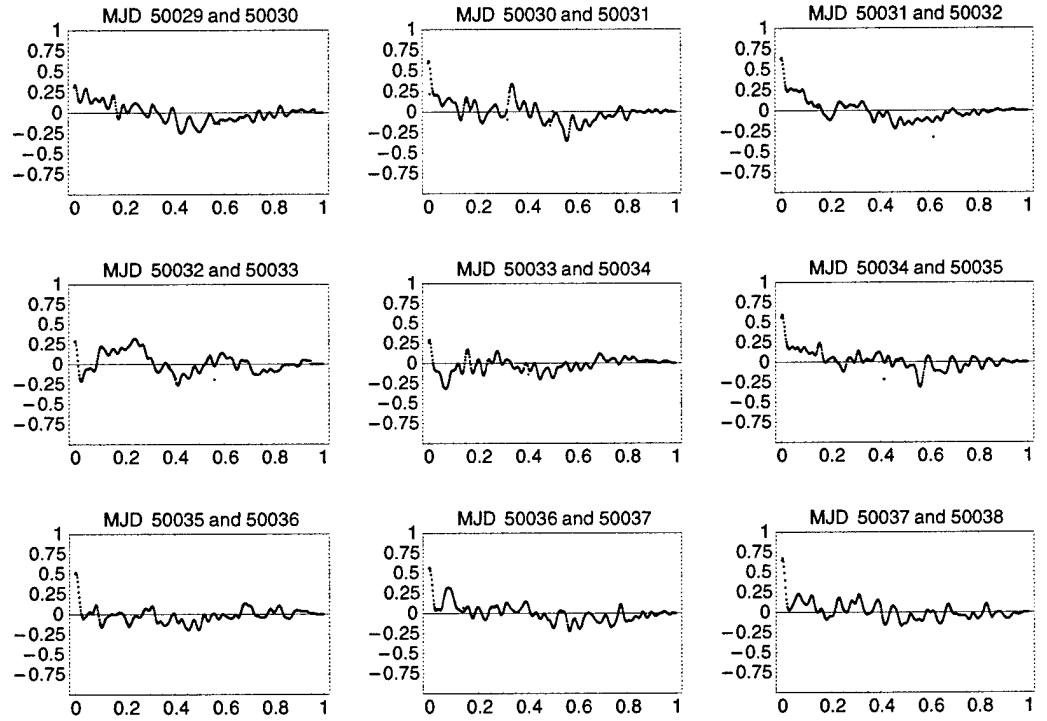


Figure 9: Correlation functions, $g_{n,n+1}(\tau)$, for $T = 1$ day, for successive days from MJD 50029 through 50037.

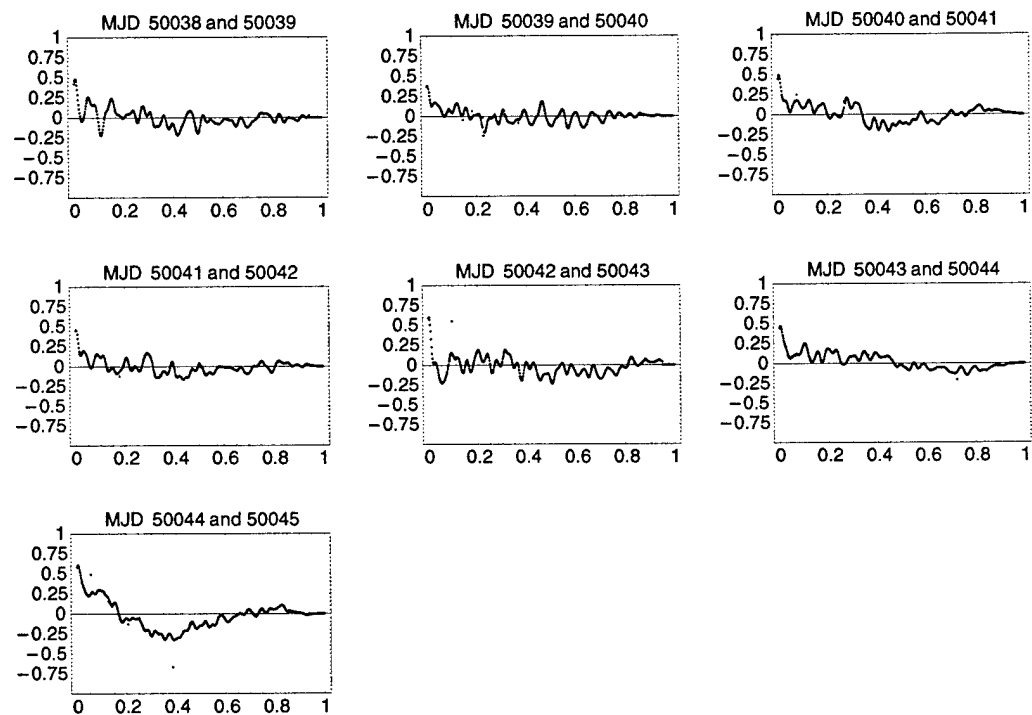


Figure 10: Correlation functions $g_{n,n+1}(\tau)$, for $T = 1$ day, for successive days starting on MJD 50038 through 50044.

that a finite sample of 25 successive-day correlation functions might show such correlations even though they do not exist; however, such an event is highly improbable. Figure 11 shows the algebraic average,

$$\langle g_{n,n+1}(\tau) \rangle = \frac{1}{25} \sum_{n=1}^{25} g_{n,n+1}(\tau), \quad (11)$$

of the 25 correlation functions. This averaging smooths out some of the fluctuations that result from using a finite data set. The averaged correlation function $\langle g_{n,n+1}(\tau) \rangle$ has a peak value of 0.52 at $\tau = 0.0030 \text{ day} \cong 4.32 \text{ min}$. Since I am taking correlations between successive days, this means that the correlation peaks at $24 \text{ hr} - 4.32 \text{ min}$, which is in good agreement with the average “24-hr” GPS satellite period— $24 \text{ hr} - 4.032 \text{ min}$.*

In the computations above, I computed correlations between two successive days. The results show that there are large correlations at an approximate delay time of 23 hr 56 min, indicating that there is a periodicity in GPS time with respect to the USNO Master Clock, over a one-day time. Using the correlation technique, we can study how long these correlations persist in time. In order to do this, I computed the autocorrelation function $g_{n,n}(\tau)$ (defined in eq (9)) for a 5-day segment of data, $T = 5 \text{ days}$, from MJD $n = 50020$ to $n = 50025$. This function, $g_{n,n}(\tau)$, is shown in figure 12. Strong peaks

*The average GPS satellite period 11.9664 hr is an average of the periods of the 25 satellites; data provided by J. Toth and B. Winn of The Aerospace Corporation.

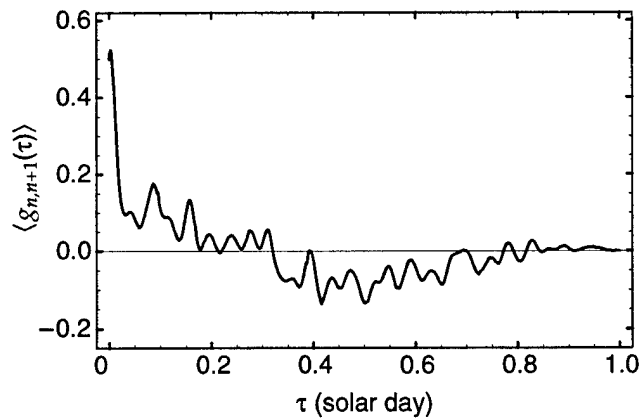


Figure 11: Algebraic average of 25 successive-day correlation functions versus delay time τ .

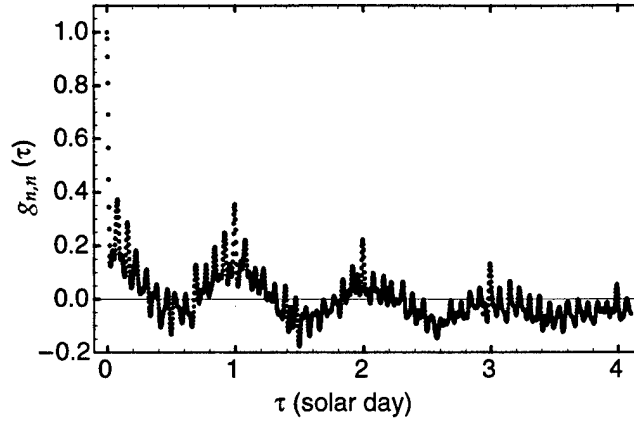


Figure 12: Autocorrelation function, $g_{n,n}(\tau)$, for $T = 5$ days, starting from MJD $n = 50020$ to 50025, is plotted versus delay time τ . Correlations persist for four successive days.

are seen at approximately 1-day intervals, indicating that the correlation at a delay time of 23 hr 56 min persists for 4 days. This “long-term memory” may be the result of long time constants present in the Kalman filter process, or systematic errors in the GPS.

At present, the source of these effects is unknown, but several sources are possible: incorrect coordinates of the receiver’s antenna phase center, errors in the GPS program code, and improper physical models being used to run the system. The strong correlation at the “24-hr” GPS satellite period suggests that some physical effects are at work.

3. Summary

The U.S. Naval Observatory Master Clock keeps the official time for DoD. Both the USNO Master Clock and the GPS are intended to keep coordinate time in the ECI (and ECEF) frames of reference [10]. Ideally, both the USNO Master Clock and the GPS are intended to keep the same time scale (modulo leap seconds); hence, if time transfer is done properly, the time difference should be zero or a random function of time with no correlations.

I have obtained time transfer data consisting of the difference between USNO Master Clock time and GPS time. A two-frequency P(Y)-code keyed receiver was used to obtain the data, and the standard two-frequency algorithm was used to correct for ionospheric delay [20].

I computed a Fourier transform of the data. Numerous sharp peaks were evident in the Fourier spectrum. A strong peak in the Fourier amplitude was seen, corresponding to approximately a 24-hr period. Two much weaker peaks were shown, corresponding to approximately 12-hr periods.

In order to determine the period near 24 hr more precisely, I performed a correlation function analysis of the data. A peak occurs in the average correlation function equal to 0.52 at a delay time of 23 hr 56 min, which corresponds to the average "24-hr" GPS satellite period. This large correlation indicates that GPS time, as seen by the U.S. Naval Observatory in time transfer, is periodic with respect to the USNO Master Clock, with period 23 hr 56 min. An autocorrelation of a 5-day segment of data shows that these correlations persisted for four successive days. The observed periodicity in GPS time indicates that there are systematic physical effects that are not adequately modelled in the GPS. Further investigations to determine the origin of these effects may lead to an improvement in overall GPS performance.

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